

Infrared dielectric metamaterials from high refractive index chalcogenides

H. N. S. Krishnamoorthy^{1,3}, G. Adamo^{1,3}, J. Yin¹, V. Savinov², N. I. Zheludev^{1,2} & C. Soci¹

High-index dielectric materials are in great demand for nanophotonic devices and applications, from ultrathin optical elements to metal-free sub-diffraction light confinement and waveguiding. Here we show that chalcogenide topological insulators are particularly apt candidates for dielectric nanophotonics architectures in the infrared spectral range, by reporting metamaterial resonances in chalcogenide crystals sustained well inside the mid-infrared, choosing Bi_2Te_3 as case study within this family of materials. Strong resonant modulation of the incident electromagnetic field is achieved thanks to the exceptionally high refractive index ranging between 7 and 8 throughout the 2–10 μm region. Analysis of the complex mode structure in the metamaterial allude to the excitation of circular surface currents which could open pathways for enhanced light-matter interaction and low-loss plasmonic configurations by coupling to the spin-polarized topological surface carriers, thereby providing new opportunities to combine dielectric, plasmonic and magnetic metamaterials in a single platform.

¹Centre for Disruptive Photonic Technologies, TPI, SPMS, Nanyang Technological University, Singapore 637371, Singapore. ²Optoelectronics Research Centre & Centre for Photonic Metamaterials, University of Southampton, London SO17 1BJ, UK. ³These authors contributed equally: H.N.S. Krishnamoorthy, G. Adamo. ✉email: harish.k@ntu.edu.sg; csoci@ntu.edu.sg

Topological insulator (TI) crystals feature time reversal symmetry-protected, highly conducting surface states characterized by Dirac dispersion and spin-momentum locking of carriers that encapsulate a semi-conducting bulk^{1–3}. They are an extremely attractive class of materials for electronic⁴, spintronic⁵ and, more recently, photonic^{6,7} applications, where coupling of light to the topologically protected surface carriers may lead to propagating surface plasmon polaritons with very little scattering and other exotic phenomena, such as hybridization of spin and surface plasmons⁸. This has motivated extensive studies of electromagnetic properties of TI chalcogenide crystals over a broad range of frequencies from THz to UV^{9–15}. Interaction of electromagnetic waves with the topological surface states can be enhanced by suitably structuring the TI crystals with subwavelength units, such as resonant metamolecules, giving rise to absorption and localization of the electromagnetic field^{9,16}. TI metamaterials have been realized at THz and UV–visible frequencies, where the material response is more plasmonic due to free surface carriers and bound bulk carriers, respectively^{9,16–18}. However, there have been hardly any studies on resonant TI structures at intermediate near-infrared and mid-infrared frequencies, where the compositionally tunable refractive index is extremely high¹⁹, and optical conductivity from charge carriers in topological surface states becomes significant²⁰. Within the family of chalcogenide crystals, we select Bi₂Te₃ to demonstrate dielectric metamaterial structures in the technologically important near to mid-infrared frequency window. Bi₂Te₃ has a refractive index between 7 and 8 over the 2–10 μm spectral range, which is much larger than typical infrared dielectric materials, such as Si ($n = 3.44$)²¹, Ge ($n = 4.07$)²², PbTe ($n = 5.61$)²³, and GST ($n = 6–7.2$)²⁴.

Here we show that the high refractive index of the chalcogenide crystals can be used to generate strong infrared resonances and associated complex mode structures with surface circular currents, opening up new opportunities to couple light with spin-polarized topological surface state carriers.

Results

The high infrared refractive index of TIs. We started by carrying out first-principles calculations and spectroscopic measurements of single crystal Bi₂Te₃. Density functional theory (DFT) calculations based on the local density approximation (LDA) were employed to study the electronic band structures and optical properties of the rhombohedral-phase Bi₂Te₃ (see the crystal structure in Fig. 1a) using the Quantum ESPRESSO (QE) package²⁵. Experimental lattice parameters of bulk Bi₂Te₃^{26,27} were used to build the initial structure, and ground state geometry of the Bi₂Te₃ crystals was obtained by the total energy minimization method upon relaxing their crystal framework and atomic coordinates. We calculated the band structure of Bi₂Te₃ with and without spin-orbit coupling (SOC). Due to the presence of heavy elements such as Bi and Te, relativistic effects and SOC have significant impact on the band structure. Without SOC, the bands have the typical parabolic dispersion with a direct gap at the Γ point. However, the presence of SOC leads to band-inversion at the Γ point and a topologically non-trivial gap is induced (Fig. 1b top panel) with the Dirac dispersion appearing in the case of a thin slab (Fig. 1b bottom panel). The optical response was calculated by employing the Bethe–Salpeter equations (BSE) method with the YAMBO code, using ground-state wavefunctions from QE package^{28,29}. The imaginary part of the permittivity was determined by evaluating direct electronic transitions between occupied and higher-energy unoccupied electronic states and the real part was obtained by employing Kramers–Kronig transformation on the imaginary part. Additional details on optical response calculations can be found in ref. ²⁰. Figure 2a shows the real and

imaginary parts of the calculated complex refractive index ($\tilde{n} = n + ik$). Its dispersion is characterized by contrasting behaviors in two regions of the spectrum: (i) the short wavelength (0.25–1.50 μm) region with strong absorption (high k), resulting from interband transitions in the bulk, and (ii) the long wavelength (~2–16 μm) region featuring strong polarizability (high n). Correspondingly, the permittivity is negative throughout the 0.25–0.85 μm region ($n < k$)^{9,20}, where the material is plasmonic. The typical concentration of surface carriers in these materials (10^{12} – 10^{13} cm⁻²) results in their contribution to the optical conductivity becoming significant in the long wavelength region above 6 μm²⁰, as manifested by a small decrease of the refractive index.

Spectroscopic studies were carried out on exfoliated films of TI single crystal samples of Bi₂Te₃ with thickness ranging from 10 to 100 μm. We characterized optical properties over a broad spectral range, from the UV to the mid-infrared by means of variable angle ellipsometry (in the UV to near-infrared range) and near-infrared to mid-infrared reflection measurements, from which we extracted the experimental complex refractive index dispersion (Fig. 2b). The experimental dispersion compares well with the calculated one, showing absorption and negative permittivity at the shorter wavelengths, and strong dielectric behavior in the infrared.

The deviation from the calculated dispersion at longer wavelengths is due to a sharp decrease of the refractive index and an increase of the extinction coefficient induced by free bulk carriers from intrinsic doping. This effectively creates a third region of high refractive index and low losses between 7 and 10 μm for the crystals in our hands. The dashed lines in Fig. 2b show the experimental optical constants after removing the bulk Drude contribution for comparison with the DFT calculation results, showing fairly good agreement between the two once the intrinsic doping contribution is excluded.

Broadly tunable infrared dielectric metamaterials. Optical materials with high refractive index and low losses such as Bi₂Te₃ in the mid-infrared are in great demand for dielectric metamaterials, as they can produce strong mode confinement and narrow resonances for small form factor devices³⁰. On this premise, we fabricated infrared nanoslit arrays via focused ion beam milling on the surface of exfoliated Bi₂Te₃ crystals. The slit array geometry was chosen for geometrical simplicity and designed to have pronounced resonances across the entire infrared spectrum by varying the nanoslit length (L) from 1.0 to 4.3 μm. Representative top-view and cross-sectional SEM images of the fabricated nanoslit arrays are shown in Fig. 3a–d. Infrared microscopy reflection measurements were carried out with incident electric field vector polarized both parallel ($R_{E\parallel}$) and perpendicular ($R_{E\perp}$) to the length of the slits. The slit arrays show resonant response only when excited with perpendicular polarization. To highlight these resonances, we plot the reflection spectra in a differential form, $(R_{E\perp} - R_{E\parallel})/R_{E\parallel}$, for the various slit lengths (Fig. 3e). Two distinct resonances are clearly observed in the plot, the fundamental resonances (indicated by \blacktriangledown) and, the second-order resonances (indicated by \blacktriangledown) which appear in the measured spectral region for slits longer than 1.5 μm. Figure 3f shows the variation of the resonant wavelength as a function of slit length, clearly indicating a linear spectral red-shift with scaling factors of 2 and 1 times the slit length for the fundamental and second-order resonances, respectively³¹.

Nature of resonant modes and circular surface currents. To further our understanding of the nature of the main resonant mode, we carried out finite-element method (FEM) simulations of the nanoslit array, over the corresponding wavelength interval.

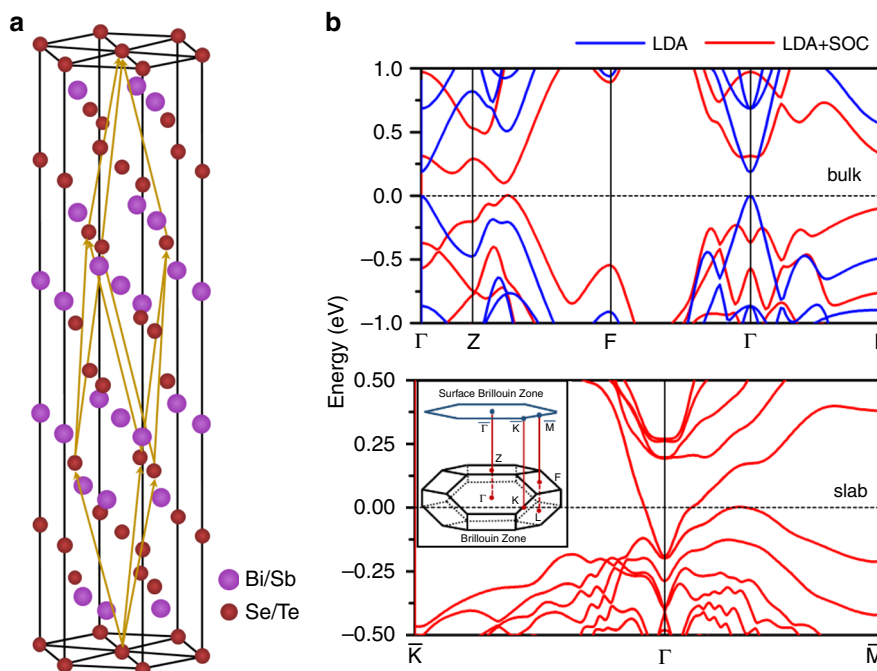


Fig. 1 Topological insulator properties of Bi_2Te_3 . **a** Crystal structure of rhombohedral-phase Bi_2Te_3 . **b** Band structure of Bi_2Te_3 bulk (top panel) and slab (bottom panel) calculated at the LDA level, with (red lines) and without (blue lines) spin-orbit coupling. Including the latter leads to band-inversion at the Γ point and gives rise to topologically protected states at the surface (bottom panel). Inset of the bottom panel shows Brillouin zone of the crystal indicating the main symmetry points.

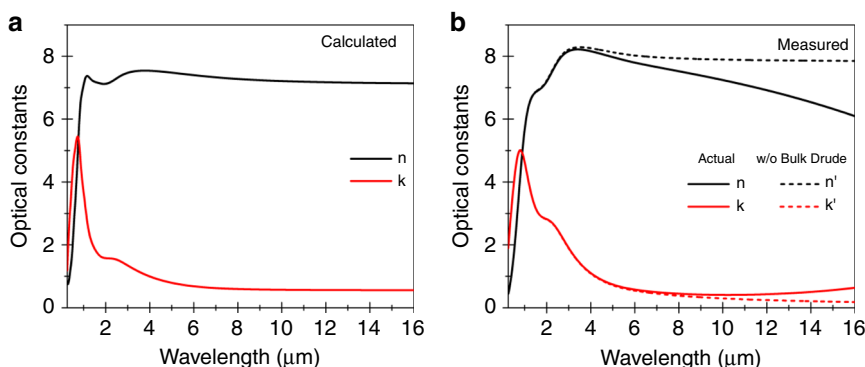


Fig. 2 Infrared refractive index of Bi_2Te_3 topological insulator crystals. Dispersion of the complex refractive index **a** calculated from first-principles technique and, **b** measured experimentally (solid lines) from ellipsometry and infrared spectroscopy measurements. The dashed lines correspond to the case in which the Drude contribution from bulk free carriers is removed.

Figure 4a compares the experimental and simulated differential reflection spectra of the longest slit ($L = 4.3 \mu\text{m}$) arrays, showing good agreement between the two. Corresponding maps of the resonant electric and magnetic fields at $8.55 \mu\text{m}$ are plotted in Fig. 4b, c for incident electric field polarization perpendicular to the slits. The field-maps for the minor mode at $4.35 \mu\text{m}$ are plotted in Fig. 4d, e. Contrary to conventional expectations, a circulating pattern for the magnetic field as displayed by the mode at $8.55 \mu\text{m}$ (Fig. 4b, c), does not necessarily imply that the resonance is a toroidal dipole. The exact multipole decomposition, shown in Fig. 4f, shows that the situation is far more nuanced: the plot of the multipole components expressed in terms of the current density (see Supplementary Eq. (12)) shows that the metamaterial excitation is not dominated by a particular mode, but is a combination of multiple modes, such as the electric

dipole (order 0), magnetic dipole and electric quadrupole (order 1), toroidal dipole, magnetic dipole, and electric quadrupole (order 2), etc., depicted schematically in Fig. 4g. Noteworthy is the fact that toroidal-like field configurations can be generated in the presence of multiple modes. The differential reflection spectrum calculated from multipole analysis is shown as dashed lines in Fig. 4a and agrees fairly well with the simulated and experimental spectra. For more details of the multipole analysis, refer to Supplementary Note 1. The multipole contribution for incident non-resonant polarization (electric field parallel to the slit length) is shown in Supplementary Fig. 2. We note that the overall reflection response of the slit arrays depends on the interference of the various multipoles based on their amplitude and phases which is rather complicated and can be illustrated with the aid of phasor diagrams (Supplementary Note 3). We also

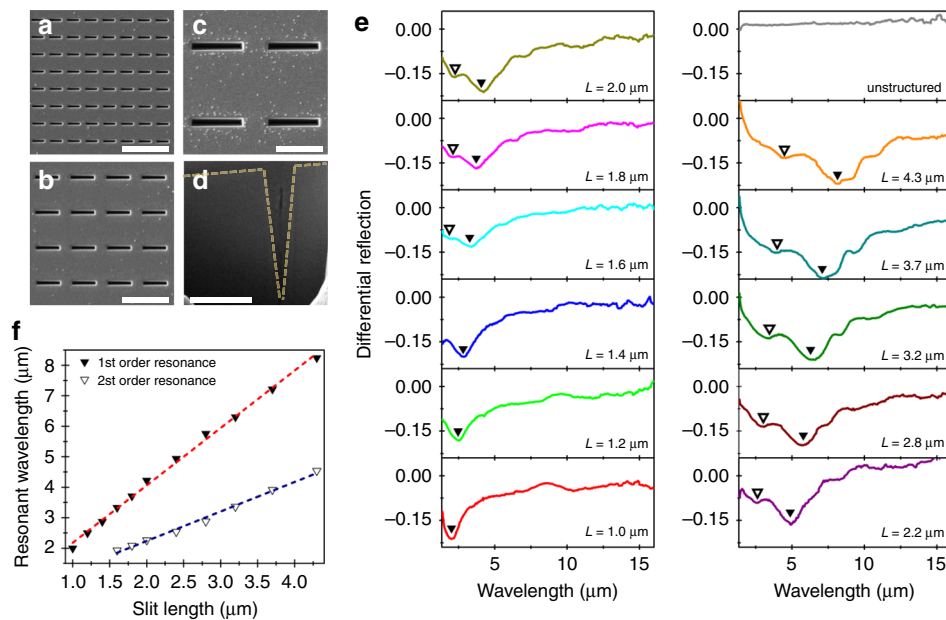


Fig. 3 Infrared resonances of nanoslit metamaterials. **a–c** SEM images of slit arrays with lengths $L = 1.0$, 2.0 , and $4.3 \mu\text{m}$, respectively. Scale bar corresponds to $4 \mu\text{m}$. **d** Cross-sectional image of slit array of length $4.3 \mu\text{m}$. Scale bar corresponds to $2 \mu\text{m}$. Dashed light yellow lines are an aid to mark the v-shaped contour of the slit. **e** Experimental infrared differential reflection spectra from slit arrays of different slit lengths from 1.0 to $4.3 \mu\text{m}$. Pronounced resonances, both first (\blacktriangledown) and second order (∇) are observed that red-shift with the slit length. **f** Resonant wavelength as a function of slit length for the fundamental and second order resonant modes in the nanoslit arrays.

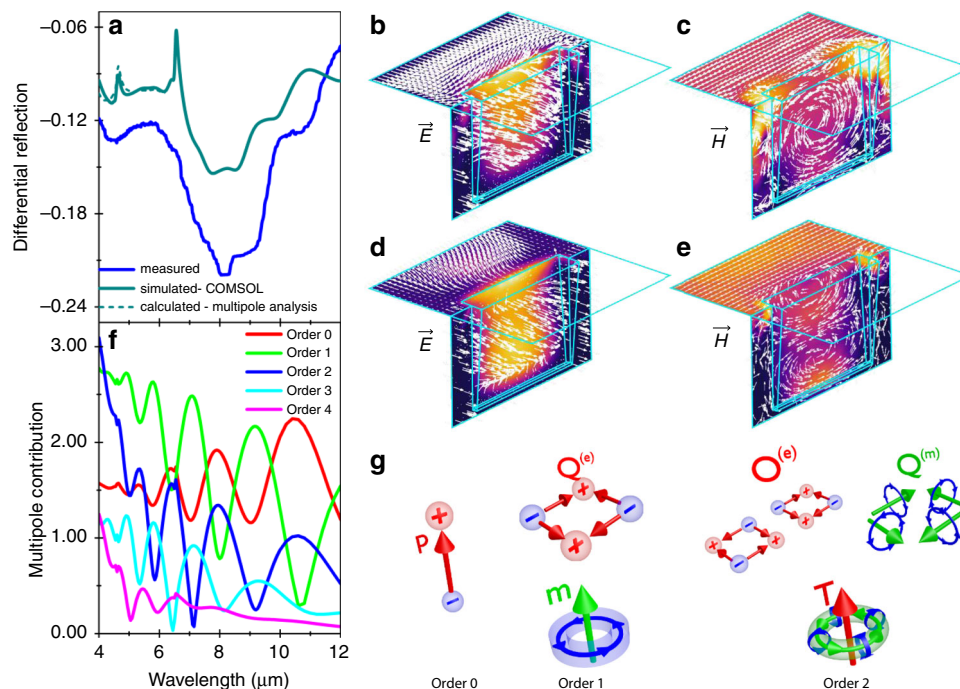


Fig. 4 Field maps and complex mode structure. **a** Experimental (solid blue), simulated (solid dark cyan) and calculated (dashed lines) differential reflection spectra of Bi_2Te_3 metamaterial slit array of length $4.30 \mu\text{m}$. The fundamental and second-order resonances at 8.55 and $4.35 \mu\text{m}$, respectively, are evident in the plot. The sharp peaks in the simulated and calculated spectra around 4.60 and $6.50 \mu\text{m}$ are due to diffraction effects. Maps of electric (\vec{E}) and magnetic (\vec{H}) fields determined by FEM simulations showing the nature of the mode at the, **b**, **c** fundamental, and **d**, **e** second-order resonances. The maps have been calculated over one half of the slit, for simulation size purposes. The cyan outline overlaid on the color maps indicates the full slit and simulation region for clarity. **f** The contribution of various multipoles to the (displacement) current excitation induced in the metamaterial by the incident wave polarized perpendicular to the slits, including electric dipole (order 0), magnetic dipole (order 1), electric quadrupole (order 1), and toroidal dipole (order 2). **g** Schematic representation of the various modes in the metamaterial. See Supplementary Note 1 for further details.

Table 1 Optical response function model.

Oscillator	Formula	Fitted parameters (cm ⁻¹)
Tauc-Lorentz	$\epsilon_2 = \frac{1}{\omega} \frac{A\omega_0 \gamma (\omega - \omega_0)^2}{(\omega^2 - \omega_0^2)^2} \Theta(\omega - \omega_0)$ $\epsilon_1 = 1 + \frac{2}{\pi} \int_0^{\infty} \frac{\omega' \epsilon_2(\omega') d\omega'}{\omega'^2 - \omega^2}$	$\omega_0 = 10,902$ $\omega_g = 1300$ $A = 7.69e^5$ $\gamma = 12,229$
Lorentz	$\epsilon(\omega) = \epsilon_{\infty} + \frac{\omega_p^2}{\omega_0^2 - \omega^2 - i\gamma\omega}$	$\epsilon_{\infty} = 1; \omega_0 = 4343.7$ $\omega_p = 19,088; \gamma = 3625.4$
Drude	$\epsilon(\omega) = \epsilon_{\infty} + \frac{\omega_p^2}{\omega_0^2 - \omega^2 - i\gamma\omega}$	$\epsilon_{\infty} = 1; \omega_0 = 0$ $\omega_p = 3147.5; \gamma = 165.6$

Parameters of the oscillators used to model the dielectric function of Bi₂Te₃ crystals throughout the optical spectral range.

note that the minor mode at 4.35 μm also shows similar behavior with circulating magnetic fields but with two vortices indicating that it is a second-order mode. Supplementary Movies 1–4 show the evolution of the electric and magnetic fields within the metamaterial over half a cycle ($0-\pi$) of the phase.

Discussion

The high-order modes supported by metamaterials^{32–34} enable strong light localization and confinement, which may be suitable for applications in nonlinear and laser optics^{35,36}, or in coupling to high-order transitions in atoms or molecules. In the context of TI metamaterials, the unique nature of the resonant fields featuring circular currents on the surface of the TI may be exploited to couple light with spin-polarized carriers and gain optical access to the topological surface states that encapsulate the material conformally³⁷. This becomes particularly relevant in the mid-IR region, where chalcogenide crystals feature a combination of high refractive index and larger contribution of topological surface states to the optical conductivity.

Overall, the TI chalcogenide crystal family is an exceptionally versatile material platform for infrared applications based on high-index, low-loss dielectric metamaterial architectures³⁸, including ultrathin flat optical elements³⁹, sub-diffraction light confinement and waveguiding⁴⁰, and nonlinear optics⁴¹. Low-loss mid-IR metamaterials are also highly sought for enhanced sensing of molecular fingerprints based on strong light confinement^{42,43}. Chalcogenide crystal metamaterials add broadband tunability of the resonances by compositional⁴⁴ and structural design to the high refractive index. Moreover, as shown here for Bi₂Te₃, the inverse geometry of nanoslits carved in high-index crystals induces higher-order complex modes outside the dielectric medium, which provide additional pathways to sense changes in the surrounding environment.

In conclusion, chalcogenide TI crystals are a compelling materials platform for photonic applications in the infrared part of the spectrum. We have shown that Bi₂Te₃ exhibits a strong polarizability with refractive index that exceeds 7 in the 2–10 μm range, larger than conventional dielectric materials that aids strong nanostructure resonances sustained deep into the mid-infrared. The exceptionally high index facilitates formation of circular currents at the surface of the material that may potentially be coupled to spin-polarized topological surface states. This opens the path to new infrared metamaterials combining dielectric, plasmonic, and magnetic properties for applications including molecular fingerprinting, environmental sensing, and integrated mid-IR photonics.

Methods

Optical characterization. Bi₂Te₃ crystals were purchased commercially from 2D Semiconductors Inc. The infrared reflection/transmission spectra of the

unstructured TIs crystals and microscopy reflection spectra of structured Bi₂Te₃ were measured using a Bruker Hyperion microscope coupled to a Bruker Vertex 80v spectrometer. Spectroscopic ellipsometry data were collected using a J.A. Woollam VASE ellipsometer in the 250–1650 nm spectral range over three angles of incidence (30°, 50°, and 70°), and analyzed using the CompleteEASE ellipsometry data analysis program. The refractive index values in the infrared were determined using the experimentally measured near-to-mid-infrared reflection spectrum and the dielectric constants in the UV–near infrared range measured from ellipsometry. This analysis was carried out using the ReFIT program wherein a combination of Tauc–Lorentz, Lorentz, and Drude oscillators were used to model simultaneously, the ellipsometric dielectric constants as well as the infrared reflection spectrum, from which the refractive index of the material in the infrared spectral range was determined. In the ReFIT program, the infrared reflection response was modeled using the model ‘Reflectivity of a Film (Epsilon + Mu) on a Substrate (code = -18)’. The parameters of the oscillators describing the material optical function are shown in Table 1.

Simulations. The optical response of the nanoslit array was simulated using full-wave Maxwell equations solver COMSOL⁴⁵. The simulations were carried out for a 3D structure using perfect electric/magnetic conductor boundary conditions on the mirror symmetry plane containing the long axis of the slit for electric field polarized perpendicular/parallel to the slit. In the multipole analysis, the number of higher-order modes obtained is dependent on the position of the air–metamaterial interface relative to the $z = 0$ plane. The optimal position of this interface is chosen such that the differential reflection as well as the field induced in the metamaterial can be quantified with the minimum possible number of modes (in this case 5), and corresponds to $z = -1200$ nm for the mode contributions shown in Fig. 4f.

Data availability

Following a period of embargo, the data from this paper can be obtained from the University of Southampton ePrints research repository, <https://doi.org/10.5258/SOTON/D1252>.

Code availability

Following a period of embargo, the codes used in this paper can be obtained from the University of Southampton ePrints research repository, <https://doi.org/10.5258/SOTON/D1252>.

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Author contributions

C.S. and H.N.S.K. conceived the idea. H.N.S.K. performed all optical measurements, extracted experimental optical constants, and drafted the manuscript. G.A. fabricated the metamaterial samples and performed COMSOL simulations. J.Y. performed DFT calculations of the bandstructure and optical constants. V.S. performed multipole analysis of the metamaterial response. C.S. and N.I.Z. supervised the work. All authors discussed the results and contributed to the revision of the paper.

Competing interests

The authors declare no competing interests.

Additional information

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Correspondence and requests for materials should be addressed to H.S.K. or C.S.

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